

## Porting the AMBER forcefield to LAMMPS—massively parallel molecular dynamics simulations of DNA

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Molecular dynamics (MD) is proving to be a valuable method for analysing the structure and flexibility of DNA (1). Atomistic MD simulations of DNA are computationally very expensive and, using conventional algorithms, are limited in practice to around 10ns (2). There are important motions within DNA which are predicted to occur over longer timescales than this (Table 1).

Table 1:– Timescales for typical DNA motions.

Time scale	Main types of internal motion
Picosecond	Short living motions and oscillations of atoms.
Nanosecond	Oscillations of small groups of atoms: sugars, phosphates, bases; bending and twisting of the double helix.
Microsecond	Winding and unwinding of the double helix; opening of base pairs.
Millisecond	Dissociation of the double helix; super helicity; overall rotation.
Second	Writhing; isomerisation; division of bacteria.

### Porting of the Forcefield

LAMMPS (Large Atomic/Molecular Massively Parallel Simulator) (3) is a parallel MD code with accurate treatment of long-range electrostatic interactions (a particularly important consideration in DNA simulations) based on the PPPM (particle-particle/particle mesh) algorithm and Ewald summation to handle the periodic boundary conditions. The advantage of LAMMPS over other codes is its ability to run calculations on parallel machines with large numbers of processors without great loss of efficiency (4), increasing the size and complexity of systems one is able to study.

We have ported the AMBER forcefield (5), which is well established for the simulation of DNA dynamics, into LAMMPS. The aim of this study has been to test this porting by comparing the dynamics of the DNA dodecamer d(CTTTTGCAAAG) as predicted by LAMMPS with dynamics data from our previous extensive analysis of this sequence using AMBER (6).

Minor changes to the LAMMPS code were required to cope with AMBER's treatment of 1-4 non-bonded interactions. Static energy calculations then showed excellent agreement between AMBER and LAMMPS (table 2).

Table 2: Static energy analysis of a representative structure of d(CTTTGC AAAAG)<sub>2</sub>.

ENERGY TYPE	AMBER (Kcal/mol)	LAMMPS (Kcal/mol)
Bond	0.0239	0.0239
Angle	399.8833	399.8833
Dihedral	438.7989	438.7989
Total VDW	2532.0560	2532.2143
Total Electrostatic	-27167.3825	-27167.0726
Total Energy	-23796.6204	-23796.1522

### Temperature Control

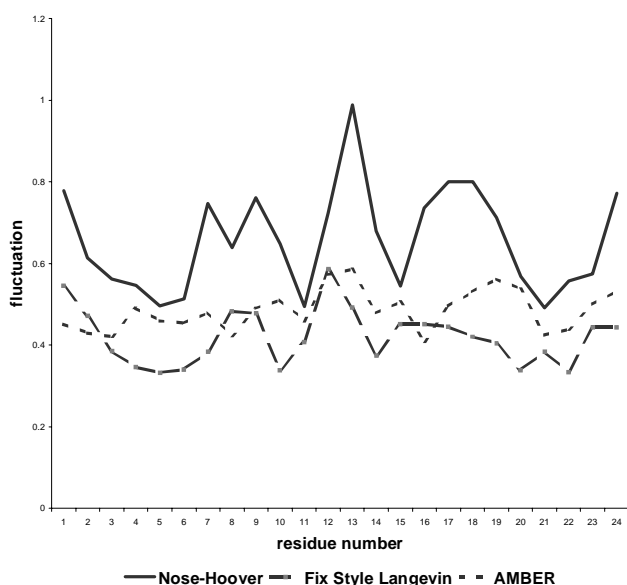
AMBER implements the Berendsen algorithm to control temperature but this is not available in LAMMPS. Two alternative temperature controls are available in LAMMPS, Langevin and Nose-Hoover. The latter approach was abandoned due to a hot solute/cold solvent problem. We have adopted the “Fix Style Langevin” temperature control where the solvent and solute temperatures are scaled separately.

Table 3: Average temperatures during 10ps runs using Nose-Hoover and Fix Style Langevin temperature couplings.

	Average Temperature (K)	Standard Deviation
Nose-Hoover Solute	320.4484	11.5720
Nose-Hoover Solvent	296.7024	5.8680
Nose-Hoover Whole system	299.7040	5.0824
Fix Style Langevin Solute	300.7899	8.9964
Fix Style Langevin Solvent	299.4550	3.9176
Fix Style Langevin Whole system	299.6238	3.5603

The hyperflexibility of the DNA, particularly at the termini of the double helix, when Nose-Hoover temperature coupling was used are apparent from the analysis of atomic fluctuations (Graph 1).

Graph 1 - Atomic fluctuation data for 10ps runs of Nose-Hoover and Fix Style Langevin temperature couplings compared to AMBER.



### Scaling and Efficiency

We are using the 816 node Cray T3E-1200E supercomputer at CSAR, the UK's national supercomputing facility, to run dynamics. Short (10ps) simulations on up to 64 processors indicate that the code is relatively efficient (equation 1) and scales well, although not linearly, in comparison to AMBER which performs much more poorly in parallel situations (see graph 2).

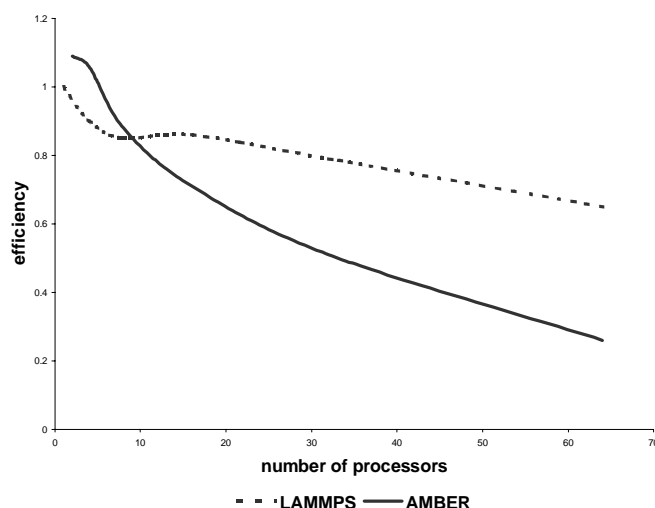
$$\text{Efficiency} = np \cdot t / l \quad (\text{Equation 1})$$

Where  $np$  = number of processors

$t$  = time taken (s)

$l$  = length of simulation (fs)

Graph 2 - Efficiency of LAMMPS compared to AMBER, normalised to the efficiency of LAMMPS on 1 processor.



### ANALYSIS OF LONGER SIMULATIONS

Three longer simulations (3ns+) have been run using 64 processors and analysis has been carried out on 2ns equilibrated portions of these. The three simulations differ only in temperature rescaling parameters, LAMMPS 1 = 0.01, LAMMPS 2 = 0.001 and LAMMPS 3 = 0.0001 (rescaling parameter in inverse time units ( $\text{fs}^{-1}$ ) therefore LAMMPS 3 has most relaxed temperature rescaling).

Three analysis techniques were used to compare the LAMMPS simulations to an AMBER simulation previously carried out:

**1. RMSD.** RMSD's have been calculated between the time-averaged structure from AMBER and corresponding time-averaged structures from the three LAMMPS simulations.

**2. ENTROPY.** Configurational entropies (7) have been calculated to obtain an overall representation of the flexibility of these systems compared to AMBER.

**3. PRINCIPAL COMPONENT ANALYSIS (PCA).** PCA has been used to identify the major modes of motion in each trajectory. The similarity between these modes in the AMBER simulation and in each LAMMPS simulation was characterised by calculating the overlap of the top 10 modes (eigenvectors) (8).

Table 4 – Characteristics of LAMMPS simulations compared to the original AMBER simulation.

SYSTEM	RMSD (Å)	ENTROPY Kcal/mol	PCA OVERLAP
AMBER	n/a	693.67	n/a
LAMMPS 1	0.9483	589.79	0.4214
LAMMPS2	0.6486	637.56	0.6309
LAMMPS3	0.6297	646.17	0.6678

### Summary

The AMBER forcefield for DNA has been successfully ported into LAMMPS, this is confirmed by excellent agreement of static energies.

We have determined (some of) the optimal parameters for stable simulations of DNA in LAMMPS. We have shown that, given the same problem, LAMMPS scales better than AMBER on a T3E. We have shown that the time-averaged behaviour of the DNA is well preserved between the two simulation techniques and that dynamical characteristics are also well preserved. It is clear from this data that LAMMPS has great potential in this field. Given large scale access to the T3E at CSAR, we would expect to be able to probe hitherto unprecedented regimes of dynamical behaviour.

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### References

- (1) Sherer E.C., Harris S.A., Soliva R., Orozco M., Laughton C.A., *J.Am.Chem.Soc.*, 1999, **121**, p5981.
- (2) Cubero E., Sherer E.C., Luque F.J., Orozco M., Laughton C.A., *J.Am.Chem.Soc.*, 1999, **121**, p8653.
- (3) Plimpton S.J., 1997, LAMMPS version5, CRADA Collaboration, Sandia National Laboratory, USA.
- (4) Maillet J.-B., Lachet V., Coveney P.V., *Phys.Chem.Chem.Phys*, 1999, **1**, p5277.
- (5) Case D.A. *et al.* 1999, AMBER 6, University of California, San Francisco.
- (6) Gavathiotis E., Sharman G.J., Searle M.S., *N.A.R.* 2000, **28**, (3) p728.
- (7) Schlitter J., *Chem.Phys.Lett.*, 1993, **215**, (6) p617.
- (8) Hess B., *Phys.Rev.E*, 2000, **62**, (6) p8438.